

*Lezioni dottorali, Università degli  
Studi di Napoli "Parthenope"*

Napoli, 21–27 October 2016

# ***Nonlinearity, Tipping Points & Chaos in the Climate Sciences***

**Michael Ghil**

**Ecole Normale Supérieure, Paris, and  
University of California, Los Angeles**



*Please visit these sites for more info.*

<https://dept.atmos.ucla.edu/tcd>

<http://www.environnement.ens.fr/>

# Overall Outline

- **Lecture I: EBMs<sup>(+)</sup>, paleoclimate & “tipping points”**
- **Lecture II: The wind-driven ocean circulation**
- **Lecture III: Advanced spectral methods—SSA<sup>(±)</sup> *et al.***
- **Lecture IV: Nonlinear & stochastic models—RDS<sup>(❖)</sup>**

(+) EBM = Energy balance model

(±) SSA = Singular-spectrum analysis

(❖) RDS = Random dynamical system

***Dottorato di ricerca***

***“Fenomeni e rischi ambientali”***

***Serie di quattro lezioni dal Professore M. Ghil***

**NONLINEARITY, TIPPING POINTS AND  
CHAOS IN THE CLIMATE SCIENCES**

con il seguente orario:

Venerdì 21,	ore 11-13,	Aula 6	(I piano Sud)
Lunedì 24,	ore 11-13,	Aula 17	(III piano Nord)
Mercoledì 26,	ore 15-17,	Aula 1	(I piano)
Giovedì 27,	ore 15-17,	Aula 1	(I piano)

# Motivation

- The *climate system* is highly *nonlinear and* quite *complex*.
- The system's *major components* — the atmosphere, oceans, ice sheets — *evolve* on many time and space scales.
- Its *predictive understanding* has to rely on the system's physical, chemical and biological modeling, but also on the thorough mathematical analysis of the models thus obtained: *the forest vs. the trees*.
- The *hierarchical modeling* approach allows one to give proper weight to the understanding provided by the models vs. their realism: back-and-forth between *“toy”* (conceptual) and *detailed* (“realistic”) *models*, and between *models* and *data*.
- How do we disentangle *natural variability* from *the anthropogenic forcing*: *can we & should we, or not?*

# ***Outline: Radiative equilibrium & EBMs***

## **1. Climate system & climate sensitivity**

- climate subsystems & other complexities
- climate sensitivity & IPCC results

## **2. Radiative equilibrium**

- global (0-D) radiative equilibrium
- latitude-dependent (1-D) radiative equilibrium

## **3. EBMs in 0-D (global)**

- equilibria & their linear stability
- multiple stationary climates

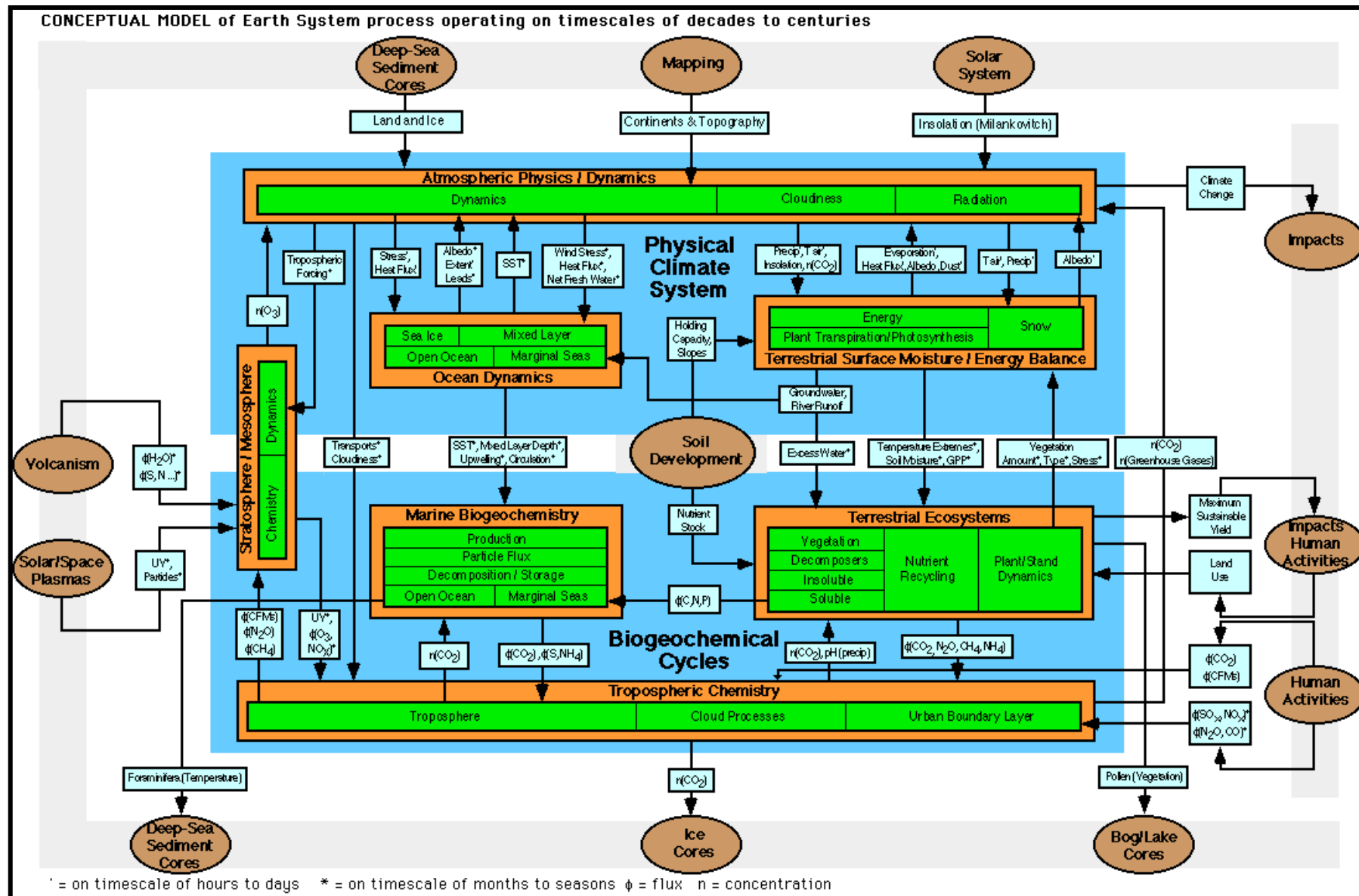
## **4. Elementary bifurcation theory in 0-D**

- saddle-node bifurcations: multiple equilibria & their stability
- nonlinear stability and variational principle
- bistability and hysteresis

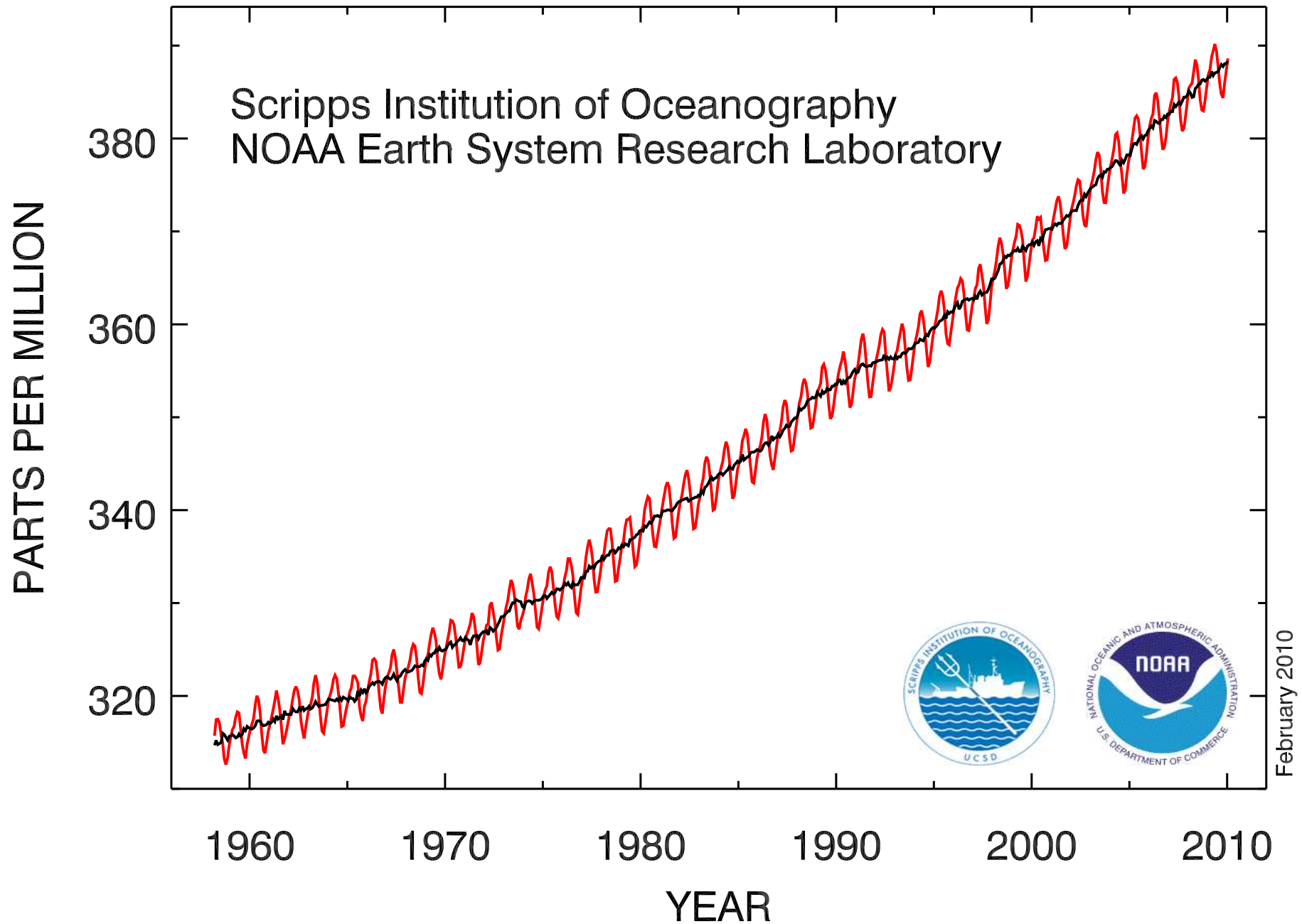
## **5. EBMs in 1-D (latitude-dependent)**

- applying all of the above
- snowball Earth

# F. Bretherton's "horrendogram" of Earth System Science



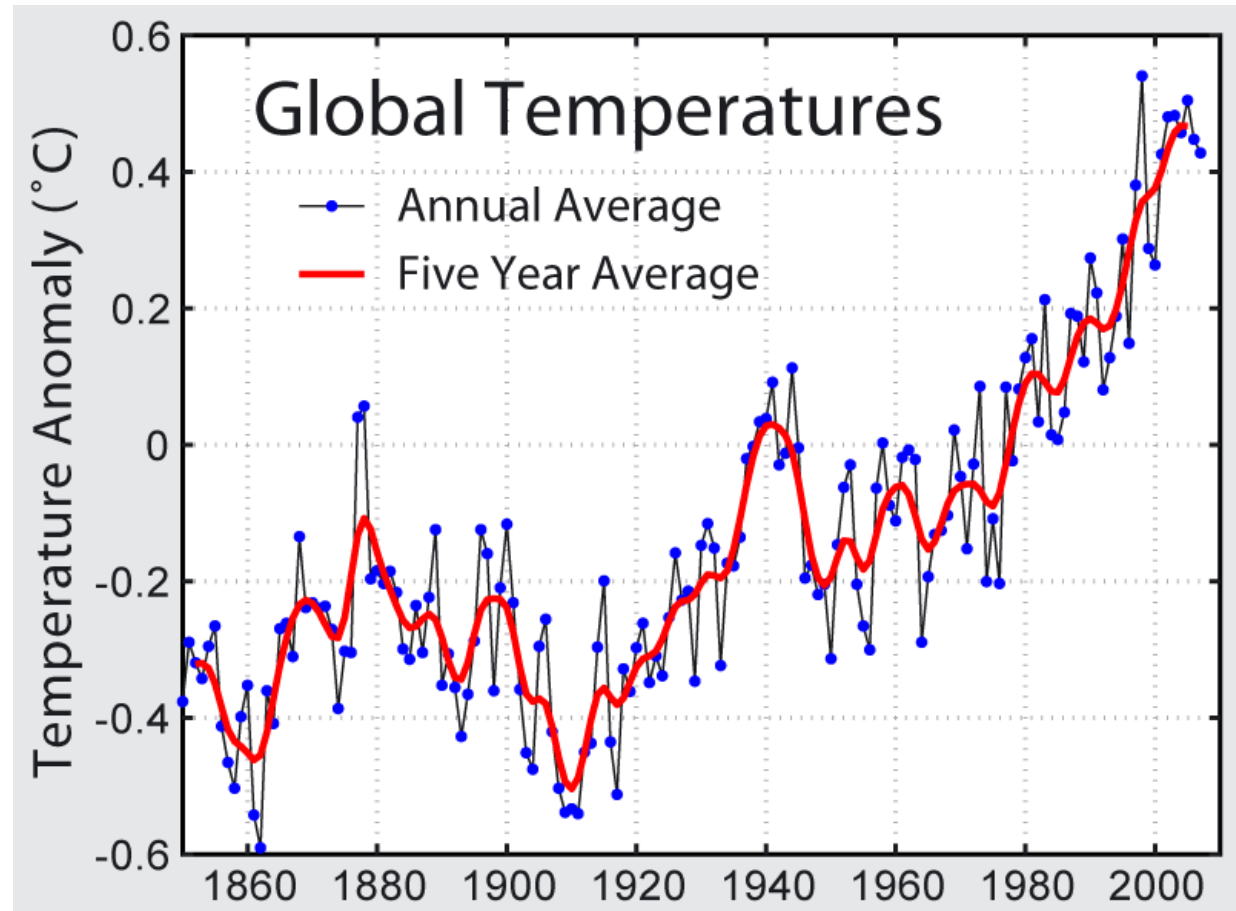
# Atmospheric CO<sub>2</sub> at Mauna Loa Observatory



# Temperatures and GHGs

Greenhouse gases (GHGs) go up,  
temperatures go up:

It's gotta do with us, at least a bit,  
doesn't it?



Wikicommons, from  
Hansen *et al.* (*PNAS*, 2006);  
see also <http://data.giss.nasa.gov/gistemp/graphs/>

# Global warming and its socio-economic impacts

Temperatures rise:

- What about impacts?
- How to adapt?

The answer, my friend, is blowing in the wind, *i.e.*, it depends on the accuracy and reliability of the forecast ...

Source : IPCC (2007),  
AR4, WGI, SPM

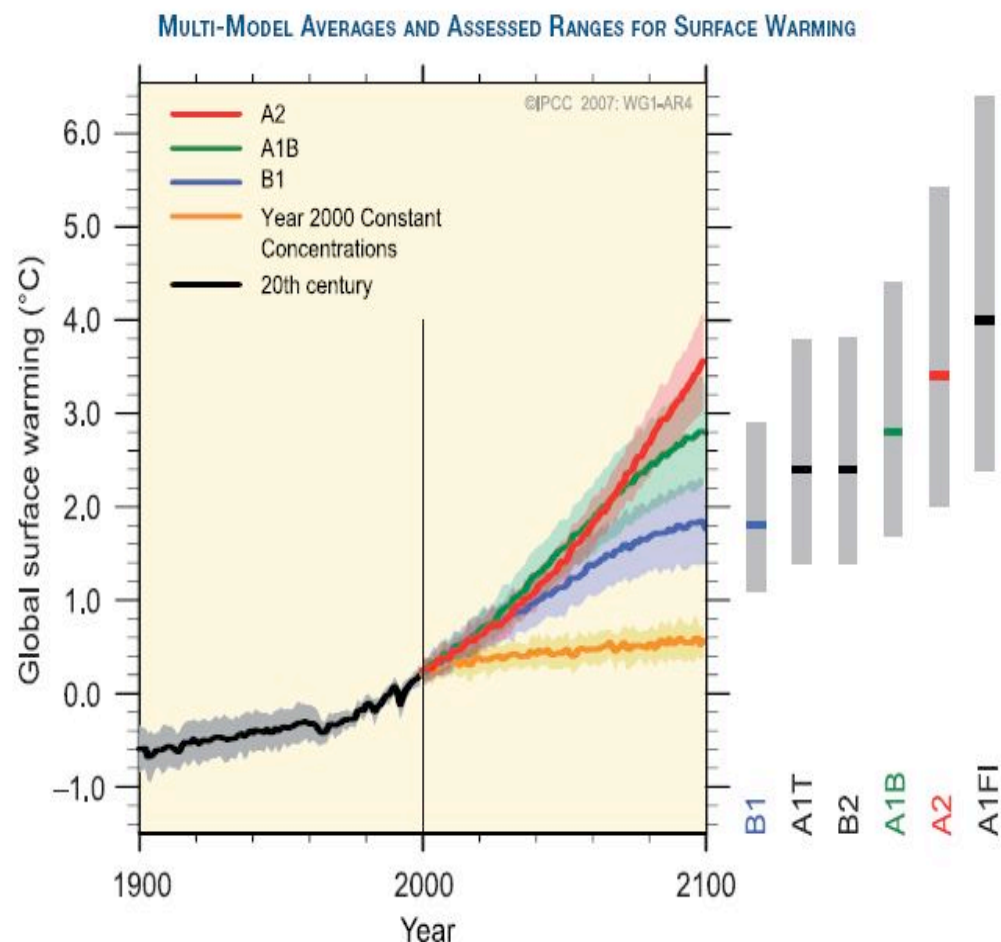


Figure SPM.5. Solid lines are multi-model global averages of surface warming (relative to 1980–1999) for the scenarios A2, A1B and B1, shown as continuations of the 20th century simulations. Shading denotes the  $\pm 1$  standard deviation range of individual model annual averages. The orange line is for the experiment where concentrations were held constant at year 2000 values. The grey bars at right indicate the best estimate (solid line within each bar) and the likely range assessed for the six SRES marker scenarios. The assessment of the best estimate and likely ranges in the grey bars includes the AOGCMs in the left part of the figure, as well as results from a hierarchy of independent models and observational constraints. (Figures 10.4 and 10.29)

# Global warming and its socio-economic impacts– II

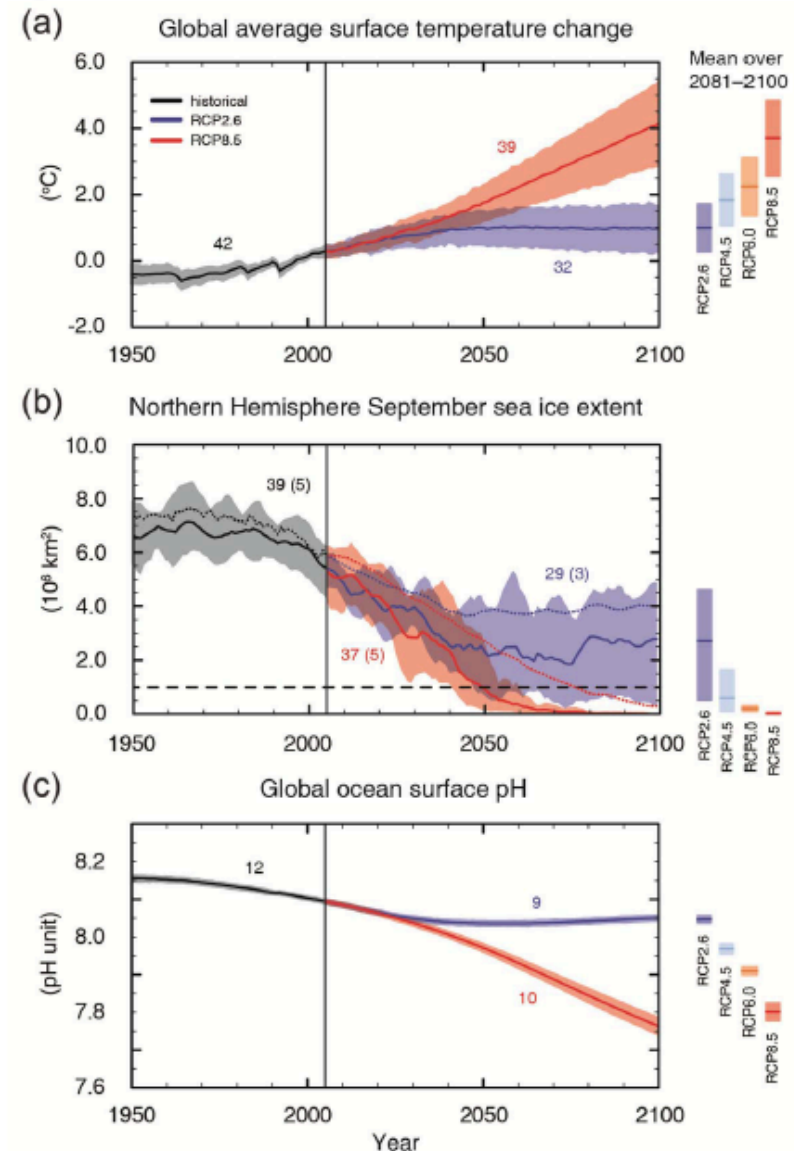
Temperatures rise:

- What about impacts?
- How to adapt?

## AR5 vs. AR4

A certain air of *déjà vu*:  
GHG “scenarios” have been replaced by “representative concentration pathways” (RCPs), more dire predictions, but the **uncertainties** remain.

Source : IPCC (2013),  
AR5, WGI, SPM



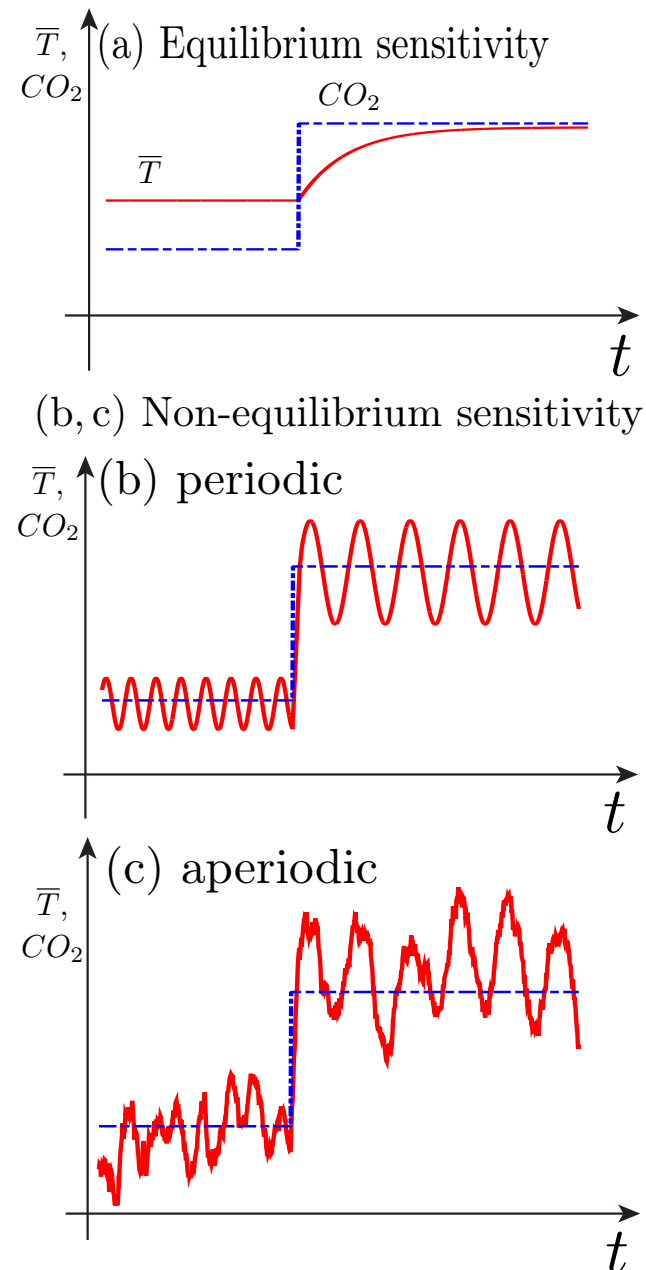
# Climate and Its Sensitivity

Let's say  $\text{CO}_2$  doubles:

How will “climate” change?

1. Climate is in **stable equilibrium** (fixed point); if so, **mean temperature** will just shift gradually to its new equilibrium value.
2. Climate is **purely periodic**; if so, **mean temperature** will (maybe) shift gradually to its new equilibrium value. But how will the **period, amplitude and phase** of the **limit cycle** change?
3. And how about some “real stuff” now: **chaotic + random**?

Ghil (in *Encycl. Global Environmental Change*, 2002)



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- equilibria & their linear stability
- multiple stationary climates

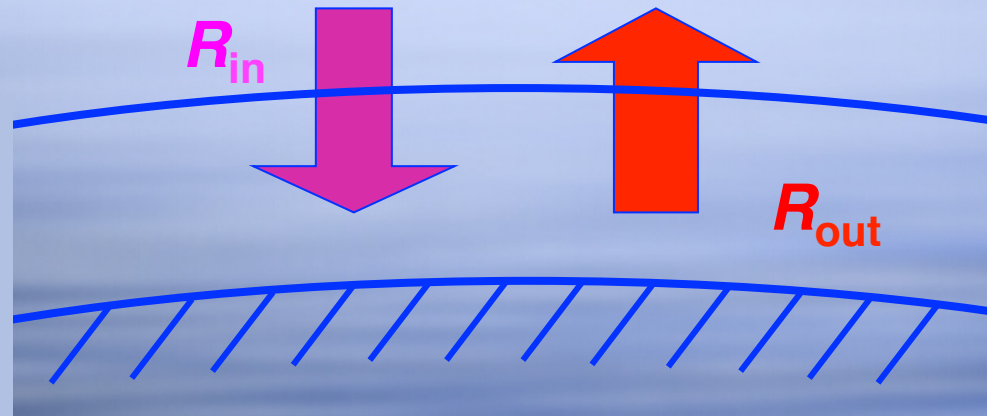
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- applying all of the above
- snowball Earth

# Radiative balance

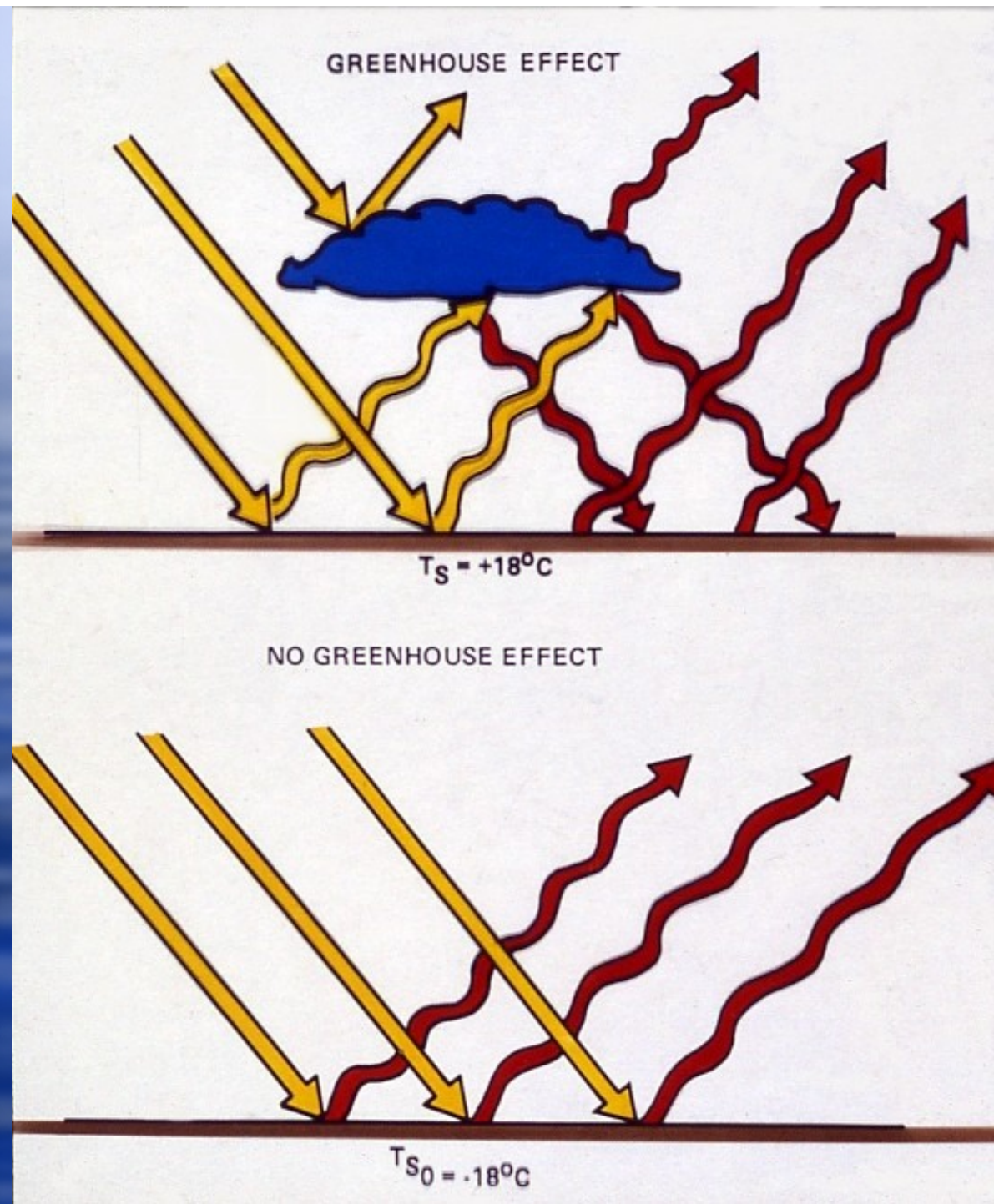


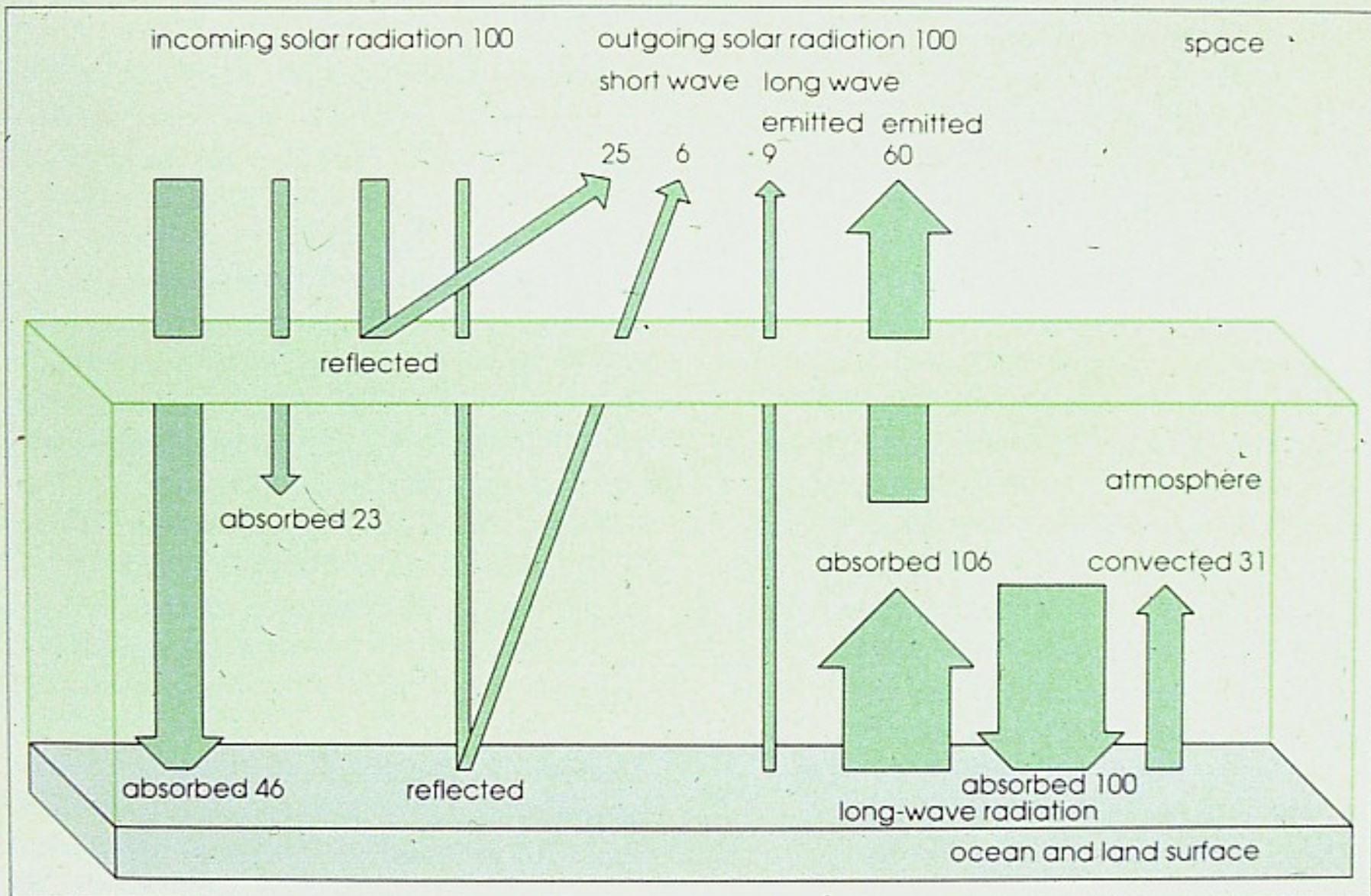
Long-term equilibrium between incident (solar, ultra-violet + visible) radiation  $R_{in}$  and outgoing (terrestrial, infrared) radiation  $R_{out}$  dominates climate.

**Refs. [1] Egyptian scribe (3000 B.C.) :**

“The Sun heats the Earth,” *Rosetta stone*, ll. 13–17.

**[2] Herodotus (484 - cca. 425 B.C.)**

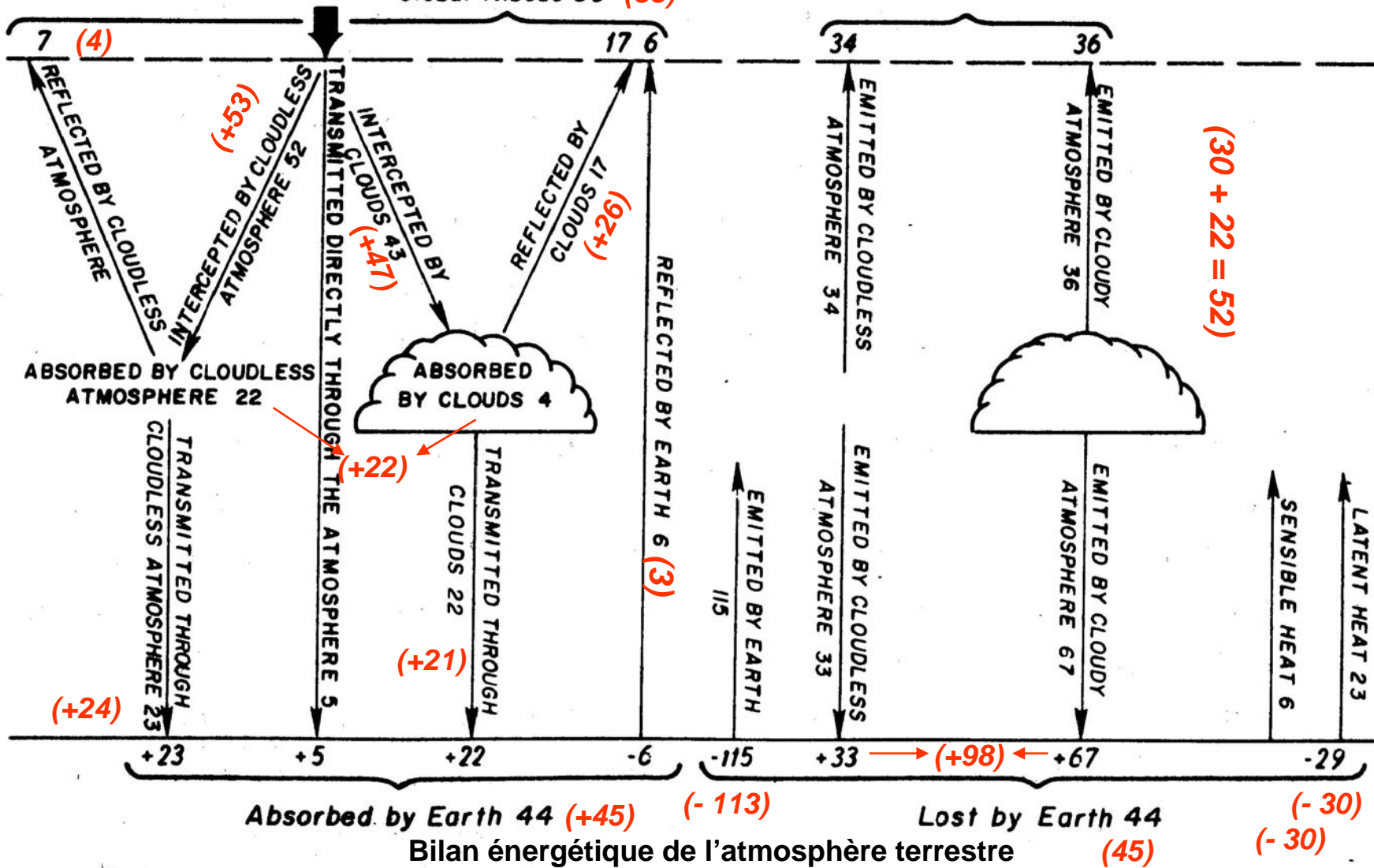




Incoming Solar Radiation 100

Infrared Heat Loss 70 (67)

Global Albedo 30 (33)

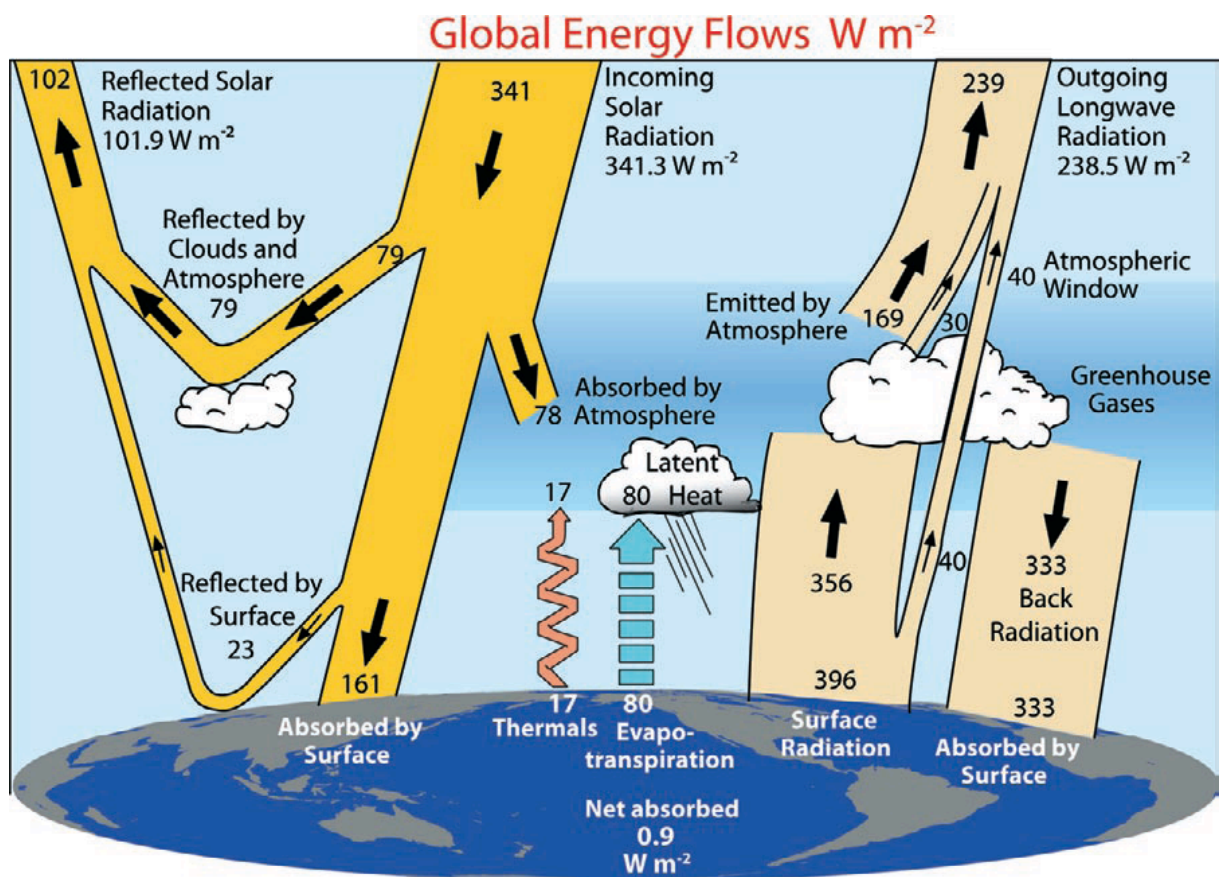


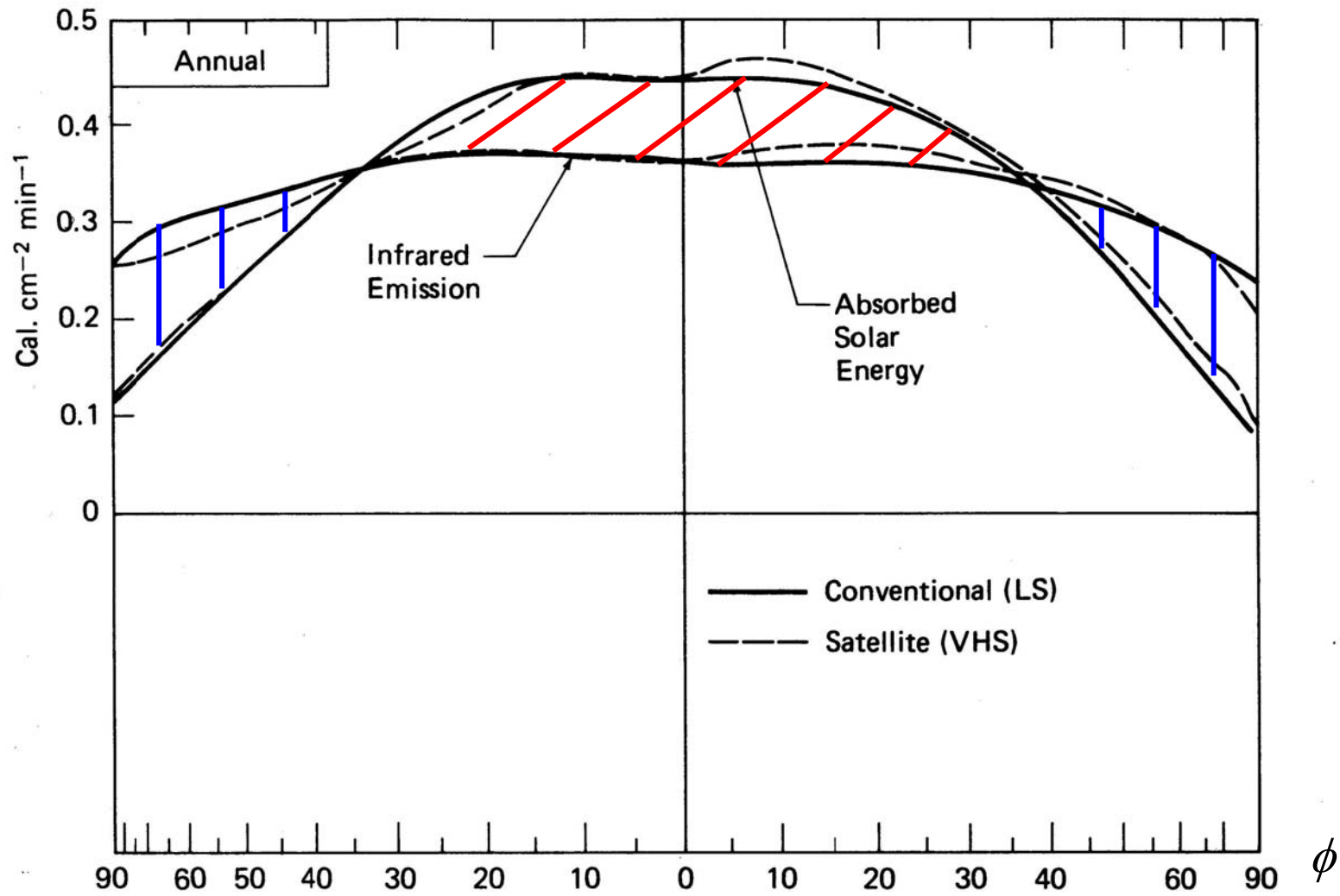
Valeurs en rouge: cf. figure précédente.

D'après Kuo-Nan Liou, 1980: *An Introduction to Atmospheric Radiation* (fig. 8.19)

# Earth's Global Energy Budget

K.E. Trenberth, J.T. Fasullo & J. Kiehl, 2009,  
*Bull. Amer. Meteorol. Soc.*, **90**(3), 311–323.





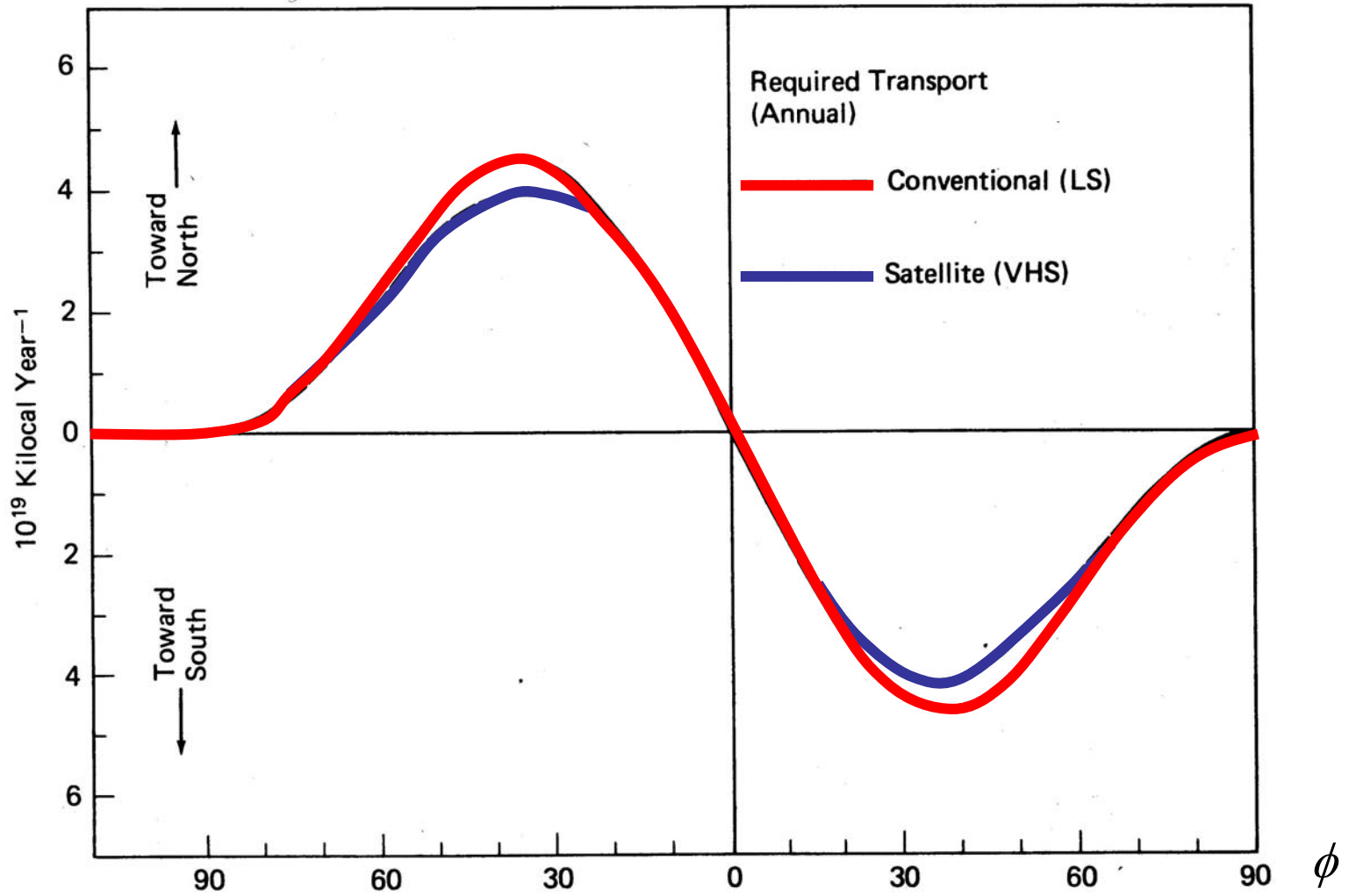
Bilan radiatif du système Terre-atmosphère en fonction de la latitude

$$R_i(\phi) - R_o(\phi) = \frac{1}{a \cdot \cos(\phi)} \frac{\partial F}{\partial \phi}$$

Avec

$a$  — le rayon de la Terre et

$F$  — le flux atmosphérique et océanique de chaleur



Transport atmosphérique et océanique d'énergie en fonction de la latitude

$$F(\phi) = a \int_0^\phi \cos(\varphi) [R_i(\varphi) - R_o(\varphi)] d\varphi$$

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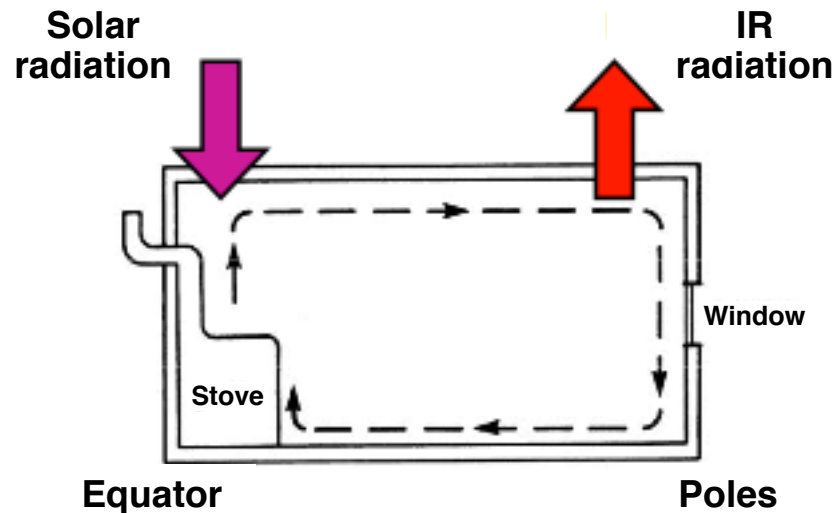
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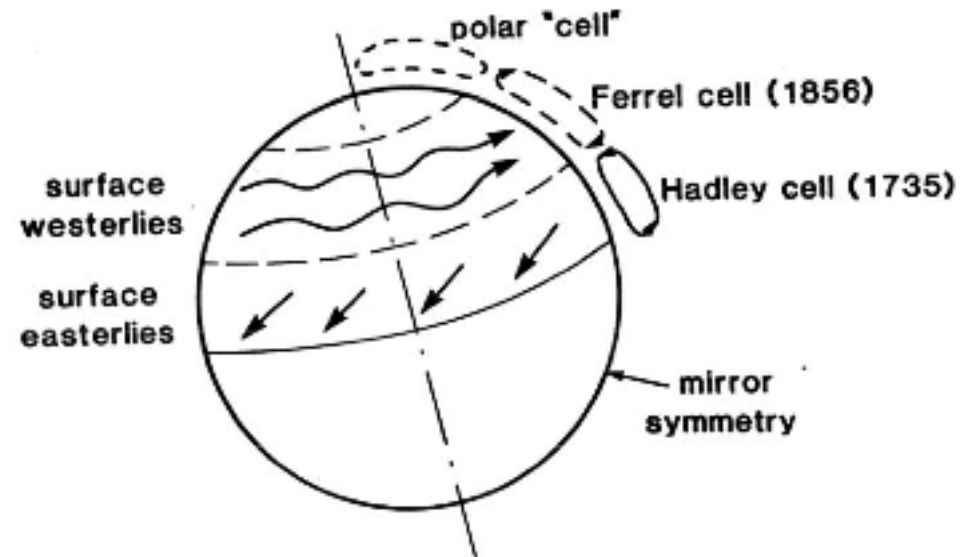
# The mean atmospheric circulation

## Direct Hadley circulation



Idealized view of the atmosphere's global circulation.\*

## Observed circulation



Schematic diagram of the atmospheric global circulation.\*

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\*From Ghil and Childress (1987), Ch. 4

# Energy-balance models (EBMs)

$$C \frac{\partial T}{\partial t} = R_i - R_o + D$$

$C$  – local calorific capacity

$T$  – local surface temperature

$R_i$  – incident solar radiation

$R_o$  – terrestrial radiation towards space

$D$  – heat redistribution ('diffusion')

## Comments:

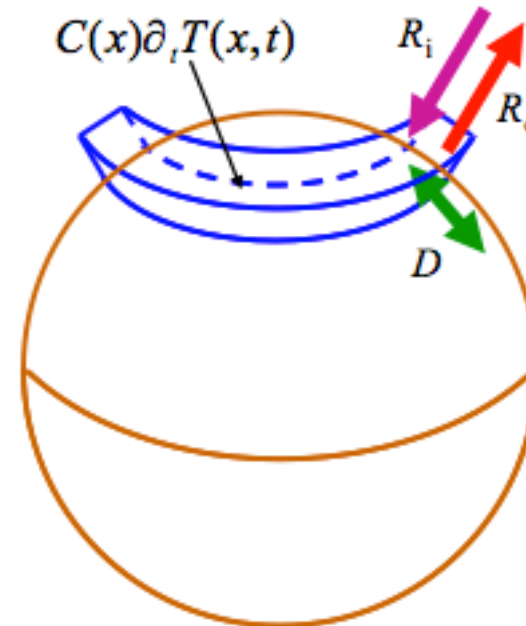
1. a space variable ( $x$ ) is maximum

2.  $C, R_i, R_o$  and  $D$  have to be calculated ("parameterized") according to  $T = T(x, t)$

3. model's main characteristic:

$$R_i = Q(x) \{1 - \alpha(x, T)\}$$

with  $\alpha$  the local albedo



## 0-D version (averaged over the globe)

$$C \frac{\partial \bar{T}}{\partial t} = R_i - R_o = Q \{1 - \alpha(\bar{T})\} - \sigma \bar{T}^4 m(\bar{T})$$

$\bar{T}$  — average surface temperature

$t$  — time (in thousands of years)

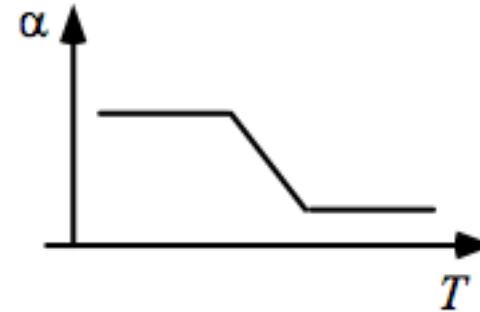
$Q$  — incident solar flux

$\alpha$  — albedo

$C$  — calorific capacity

$\sigma$  — Stefan–Boltzmann constant

$m$  — greenhouse effect factor



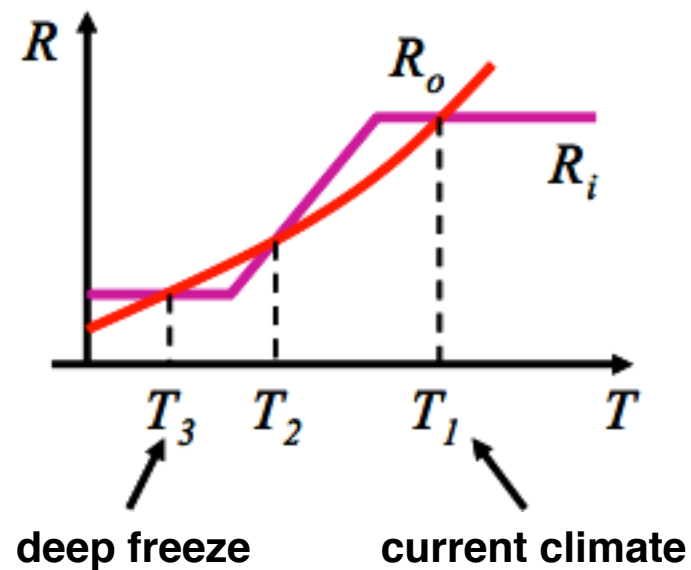
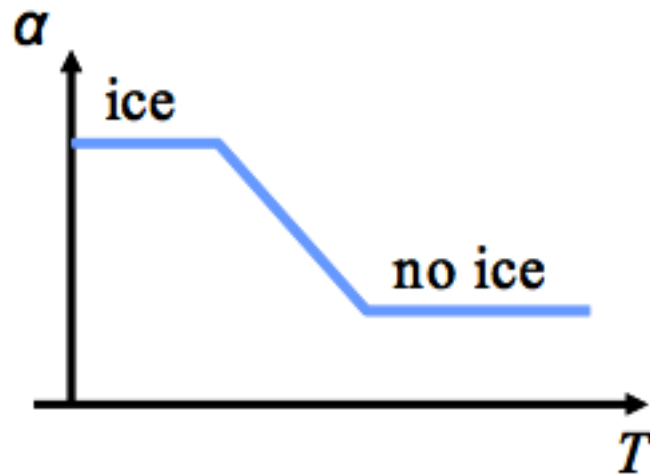
### Comments:

$\alpha$  depends on the ice and snow cover, on cloud cover, etc. (implicit variables). All is parameterized as a function of the explicit variable  $\bar{T}$ .

## 0-D EBM, I: Model solutions

We want to write  $T$  as:  $T = T(t; T_0, Q, c, \dots)$

Stationary solutions:  $Q\{1 - \alpha(T)\} - \sigma T^4 = 0$



What happens if the sun "blinks" and  $T = T_1 + \Delta T$ ?

We have to go back to the original equation, which depends on time.

## 0-D EBM, II: Stability condition

$$C\partial_t T = R_i - R_o = f(T)$$

$$R_i = Q\{1 - \alpha(T)\}$$

$$R_o = A + BT$$

We set  $T = T_j + \theta$ .

$$f(T_j) = 0,$$

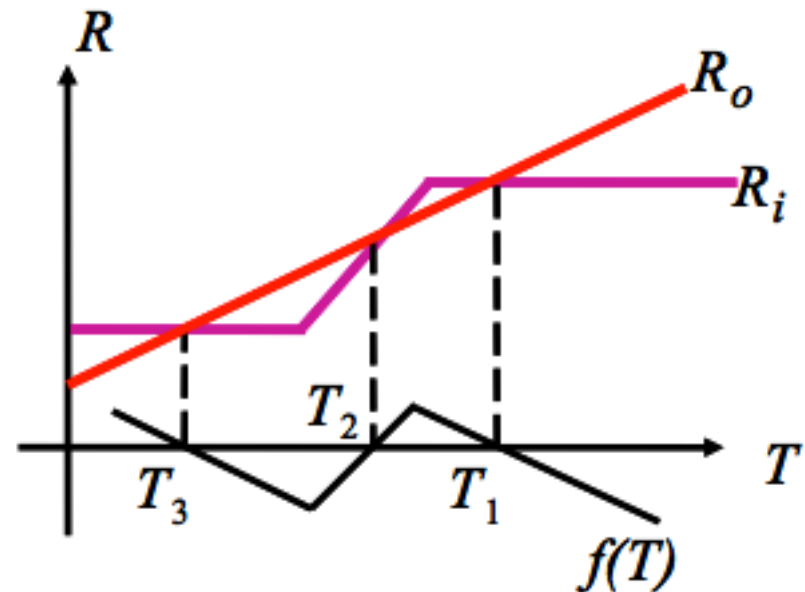
$$f(T) = f(T_j) + f'(T_j)\theta + \dots$$

Let's define  $\lambda_j \equiv f'(T_j)/c$

$$\partial_t \theta = \lambda_j \theta \Rightarrow \theta = e^{\lambda_j t} \theta_0$$

If  $\lambda_j < 0$  stable;

if  $\lambda_j > 0$   
unstable.

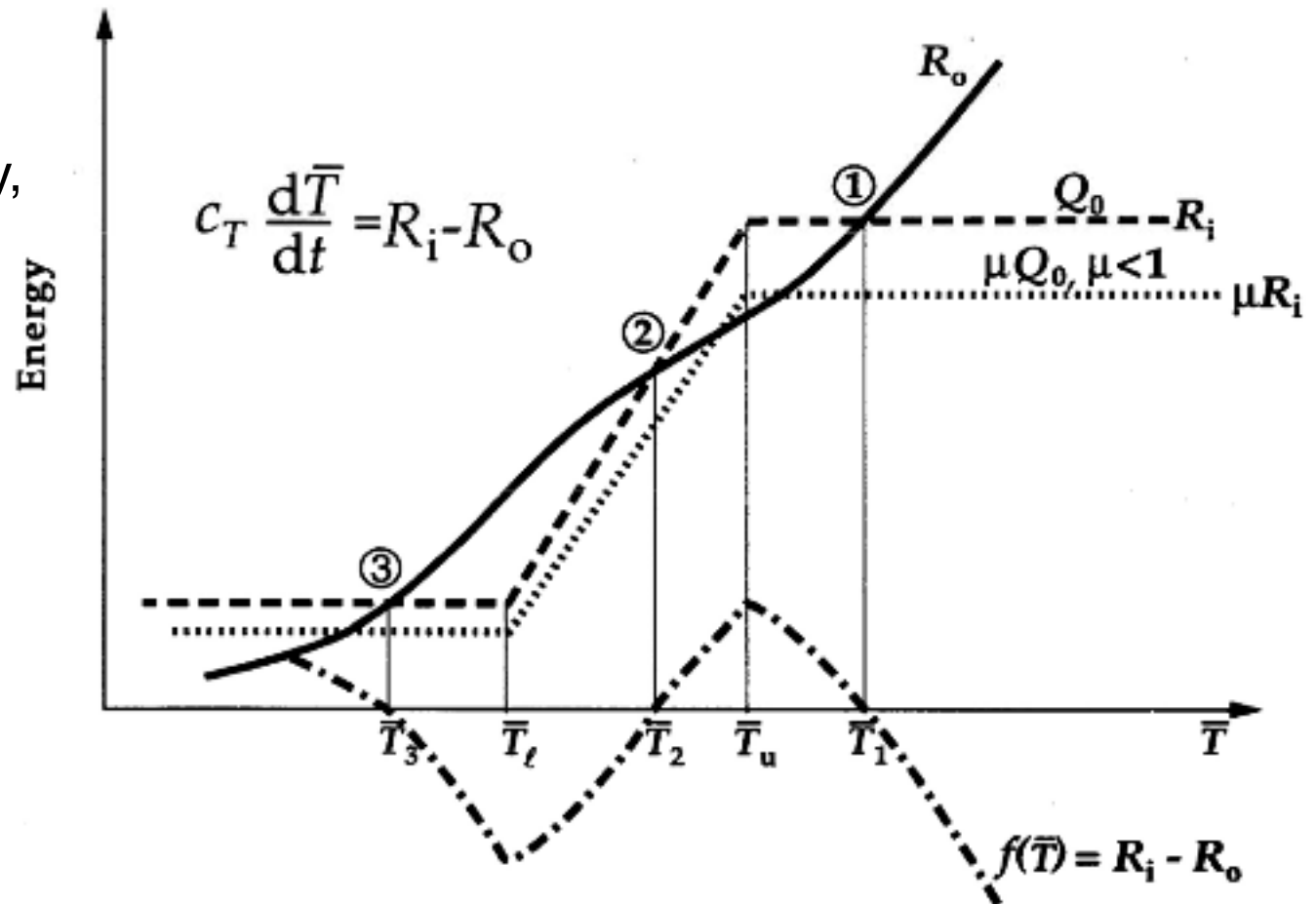


**Comment:** in the 1-D case  $\lambda_j \rightarrow \lambda_j^{(0)}$ ;  
 $\lambda_j \sim 1/c$

## 0-D EBM, III: Changes in parameters

What happens if the insolation parameter  $\mu$  changes, i.e., the “solar constant” changes? This may represent a change in solar luminosity, orbital parameters or in the optical properties of the atmosphere.

❖ The model’s three “climates” shift in value and, possibly, in number.



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# ***Outline, Tipping Points I***

## **Elementary Bifurcation Theory and Variational Principle**

### **1. Fixed Points**

- linear stability
- non-linear stability and attractor basins

### **2. Saddle-node bifurcations**

- multiple branches of stationary solutions
- linear stability

### **3. Bifurcations in 1-D**

### **4. Non-linear stability and variational principle**

- variational principle in 0-D
- variational principle in 1-D

### **5. Bistability and hysteresis**

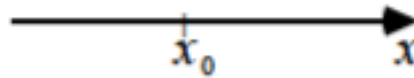
# 1. Fixed points, I

We start with a **scalar** ordinary differential equation (ODE)

$$\dot{x} = f(x; \mu)$$

depending on the parameter  $\mu$ .

**Linear stability**,  $\mu = 1$ .



$$f(x_0) = 0 \Rightarrow \dot{x} = 0 \Rightarrow x \equiv x_0 \quad - \text{Fixed point (FP)}$$

Consider an initial perturbation at  $t = 0$ :

$$x(0) = x_0 + \xi(0),$$

$$\dot{x} = \dot{x}_0 + \dot{\xi} = \dot{\xi}$$

$$= f(x_0 + \xi) = f(x_0) + f'(x_0)\xi + O(\xi^2)$$

For an infinitesimal perturbation  $\xi(0) = \xi_0$

$$\dot{\xi} = f'(x_0)\xi, \quad f'(x_0) = \lambda, \quad \dot{\xi} = \lambda\xi,$$

$$\Rightarrow \xi(t) = e^{\lambda t}\xi(0)$$

# 1. Fixed points, II

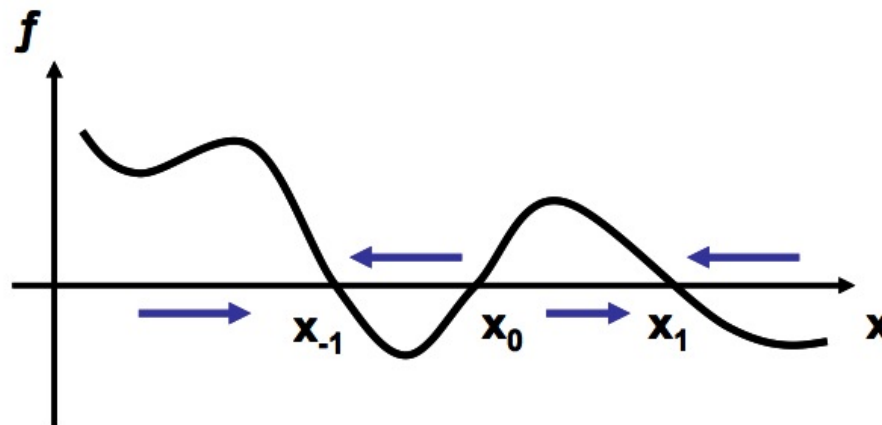
If  $\lambda < 0 \Rightarrow$  the fixed point (FP) is (linearly) **stable**

If  $\lambda > 0 \Rightarrow$  the FP is (linearly) **unstable**

If  $\lambda = 0 \Rightarrow$  the linear stability of the FP is **neutral**

*Some basic features on FPs:*

1.  $f \in C^1, f \neq 0$  on all sub-intervals: FPs are isolated (generic property)
2. Basins of attraction are open intervals (possibly semi-infinite)



## 2. Saddle-node bifurcations

How does the geometry of the solutions change when  $\mu \neq \mu_0$ , i.e. how do the number of the stability of the stationary solution change?

Let us start with the scalar case.

A simple case: the **saddle-node**

$$\dot{x} = \mu - x^2 \equiv f(x; \mu)$$

FPs:  $\mu - x^2 = 0 \quad x = \pm\sqrt{\mu}$

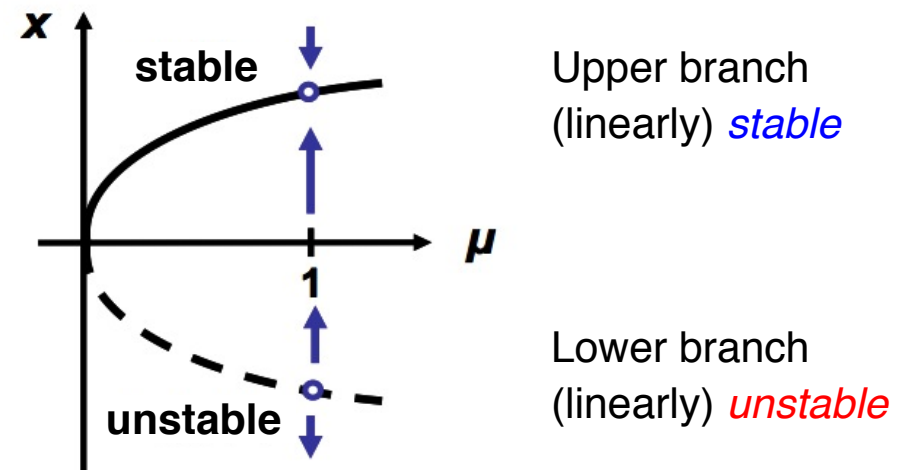
FP stability:

$$x_1 = \sqrt{\mu}, \quad x_{-1} = -\sqrt{\mu}$$

$$x(0) = x_{\pm 1} + \xi(0)$$

$$\dot{\xi} = \lambda_{\pm} \xi,$$

$$\lambda_{\pm} \equiv f'(x_{\pm 1}) = -2x_{\pm 1} = \mp 2\sqrt{\mu}$$



Let us now examine the nonlinear stability

### 3. Bifurcations in $n$ -D

We studied the scalar case ( $n = 1$ ). More generally, we have:

$$\dot{\vec{x}} = \vec{f}(\vec{x}; \mu), \quad \vec{f} \in C(\mathbb{R}^n \times \mathbb{R}), \quad \text{with } \vec{x} \in \mathbb{R}^n \text{ and } \mu \in \mathbb{R}.$$

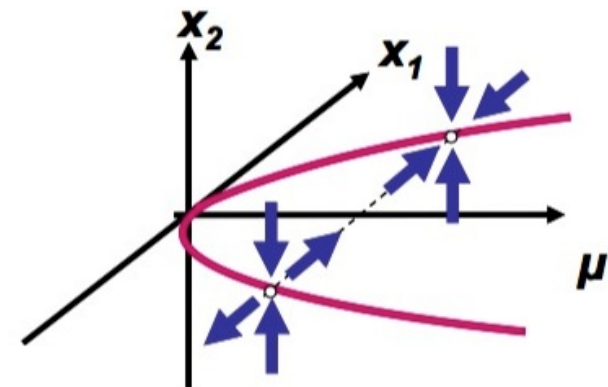
The behavior is "almost" linear in all the phase-parameter space  $\mathbb{R}^n \times \mathbb{R}$ , except in the neighborhood of a few isolated points  $(x_c, \mu_c)$ : these are **bifurcation points**, where the Jacobian matrix  $L = (\partial f_i / \partial x_j)$  is singular, i.e.  $\det L = 0$

In the case  $n = 2$ , we can reduce to the normal form :

$$\begin{aligned} \dot{x}_1 &= \mu - x_1^2 \\ \dot{x}_2 &= -\lambda x_2, \quad \lambda > 0 \end{aligned}$$

In the general case, the reduction gives :

$$\begin{aligned} \dot{x}_1 &= \mu - x_1^2 \\ \dot{x}_i &= -\lambda_i x_i, \quad \lambda_i > 0, \quad i = 2, \dots, n \end{aligned}$$



This shape explains the "saddle-node bifurcation" terminology

## 4. Non-linear stability and variational principle

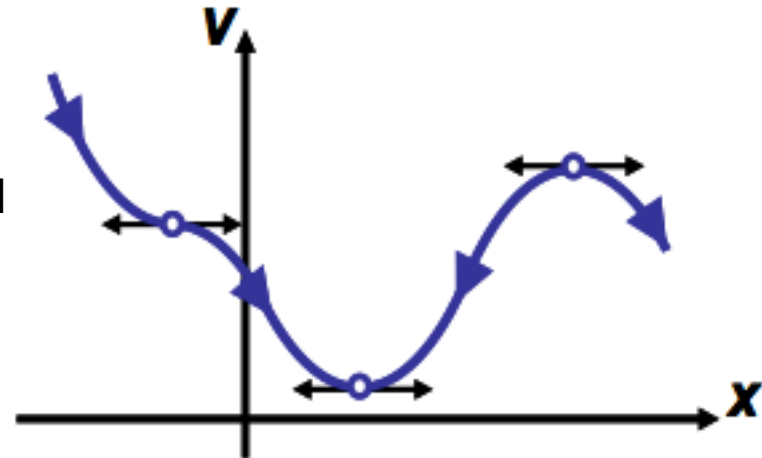
To deepen our understanding of stability, we have to examine the effect of larger perturbations.

### a) Variational principle in 0-D

$$V(x) = - \int_x f(\xi) d\xi \quad \text{— pseudo-potential}$$

$$\dot{x} = f(x) = -V'(x)$$

$$\dot{x}^2 = - \frac{\partial V}{\partial x} \frac{\partial x}{\partial t} = -\dot{V}$$



$V$  will *decrease* along the ODE's trajectory as long as :

$$\dot{x} \neq 0 \Leftrightarrow V' \neq 0$$

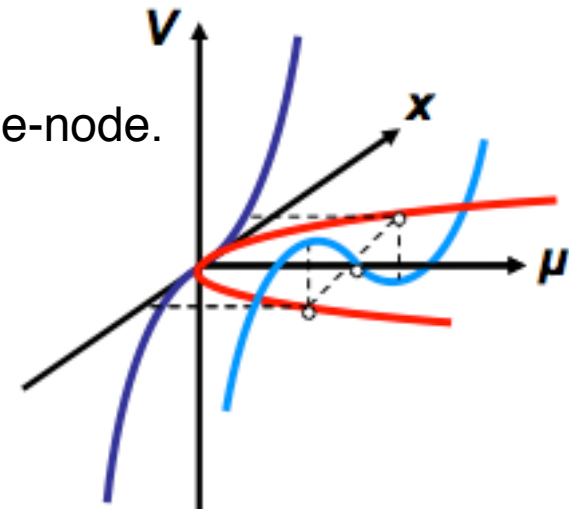
$\dot{x} = 0$  if  $V$  reaches a minimum, a maximum, or a saddle-node.

Of course, only  $V = \min$  is **stable** — *nonlinearly*.

With this result, we turn back to the saddle-node bifurcation:

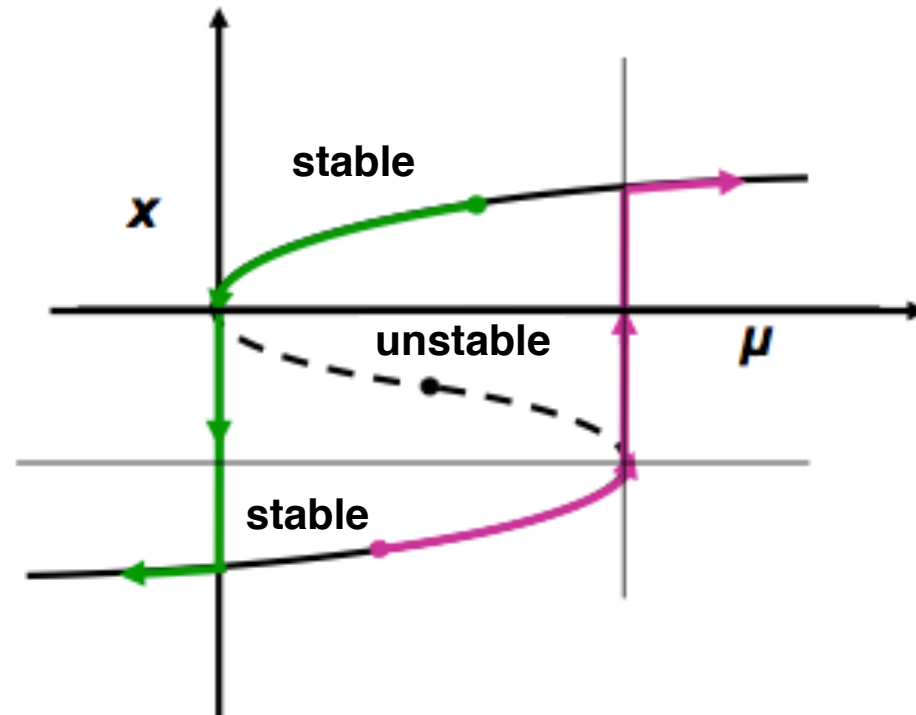
$$\dot{x} = \mu - x^2$$

$$V(x; \mu) = -\mu x + x^3/3 + c(\mu)$$



## 5. Bistability and hysteresis

The combination of two saddle-node bifurcations can create a hysteresis phenomenon (an S-shaped curve) :



$\dot{x} = \mu - x^2$  : the top-left bifurcation

$\dot{x} = (\mu - 1) + (x + \frac{1}{2})^2$  : the bottom-right bifurcation

## 1-D version ('classic' EBM)

$$C(x)T_t = R_i - R_o + D$$

$T$  — temperature

$x$  — latitudinal coordinate

$\tilde{T}(x)$  — the observed climate

Boundary conditions:  $T_x(0) = T_x(1) = 0$

$x = 0$  Pole (North)

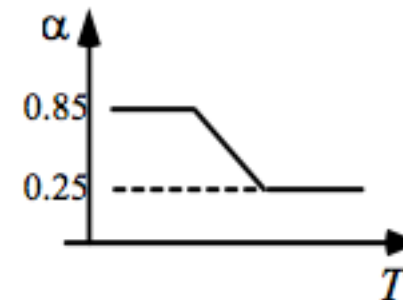
$x = 1$  Equator

$$R_i = Q(x)\{1 - \alpha\}$$

$$= Q(x)\{1 - b(x) + c_1 T\}_c$$

$$R_o = \sigma T^4 \{1 - m \tanh(c_3 T^6)\}$$

$$D = \frac{1}{\sin \frac{\pi x}{2}} \partial_x \sin \frac{\pi x}{2} \{k(x) + k_s(x)g(\tilde{T})\} T_x$$



**Questions:** 1. Stationary solutions ('climates')?

2. Stability?

3. Perturbation & bifurcation?  $Q \rightarrow \mu Q$  ( $\mu = 1$ )

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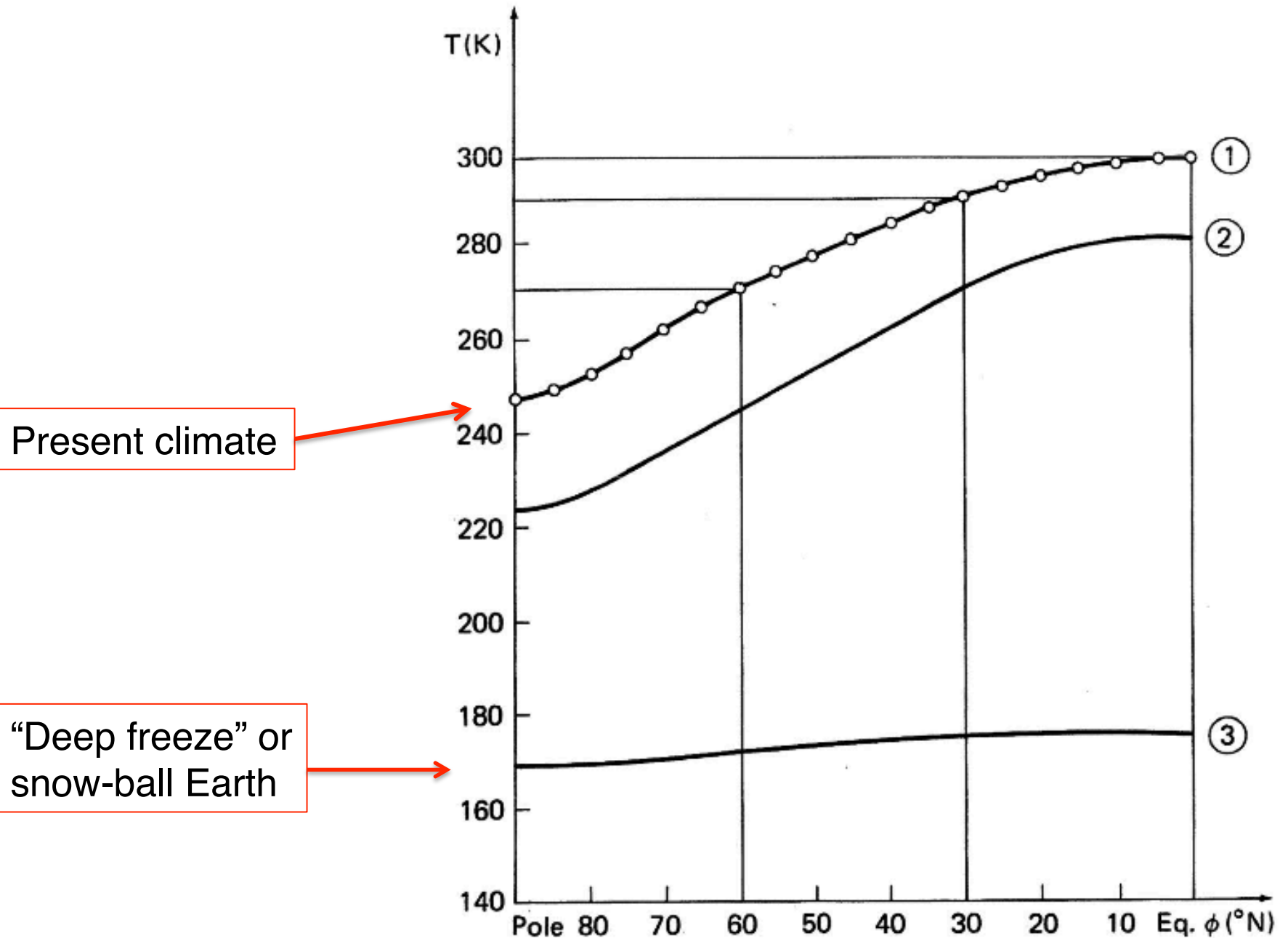
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## The three climates of the 1-D model



# 1-D EBM: Bifurcation diagram

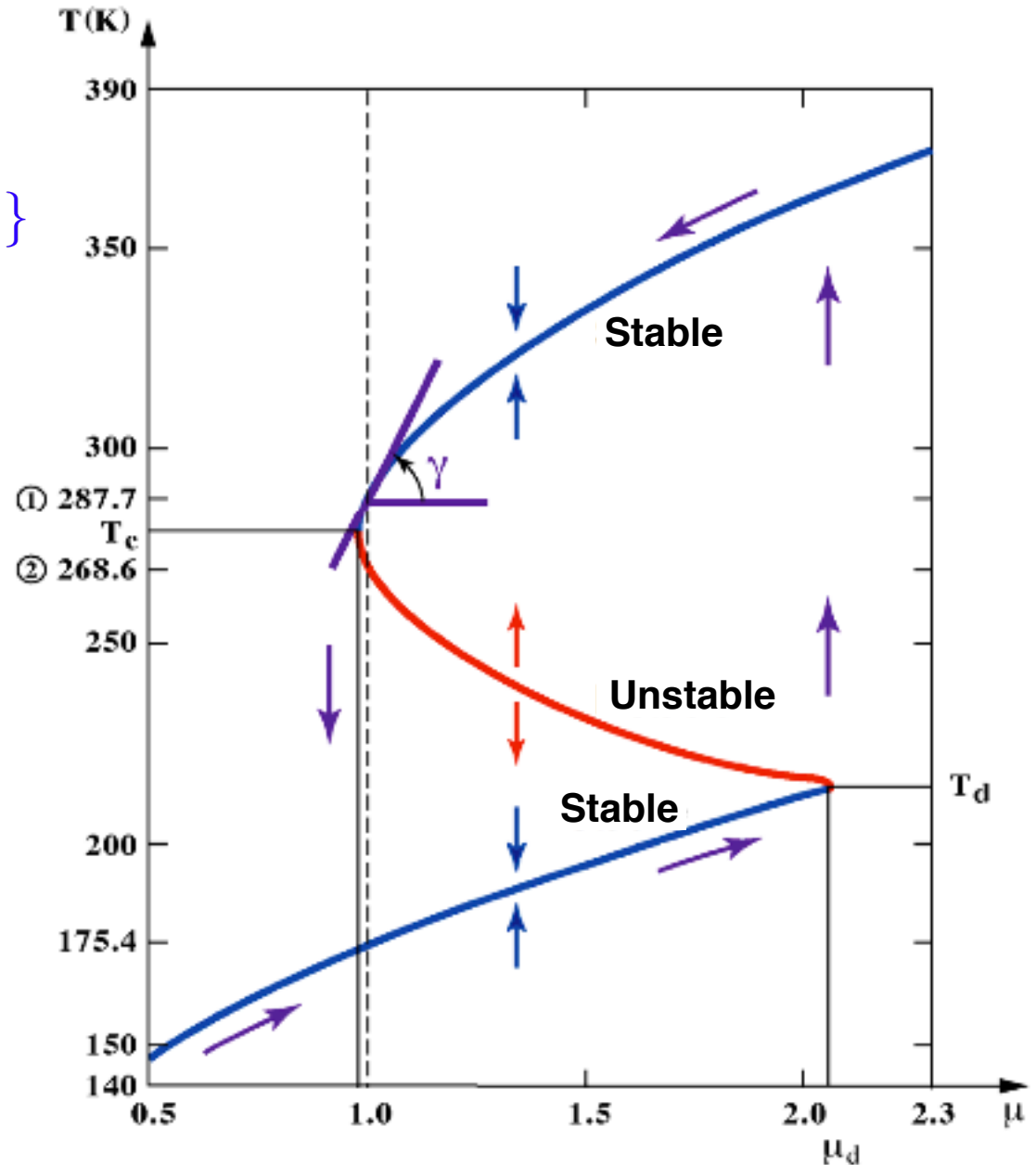
$$C(x)T_t = \{k(x, T)T_x\} \\ + \mu Q_0 \{1 - \alpha(x, T)\} \\ - g(T)\sigma T^4$$

$$T_x = 0 \text{ at } x = 0, 1$$

**Climate sensitivity:**

$$\gamma = \frac{dT}{d\mu} \cong 0.01$$

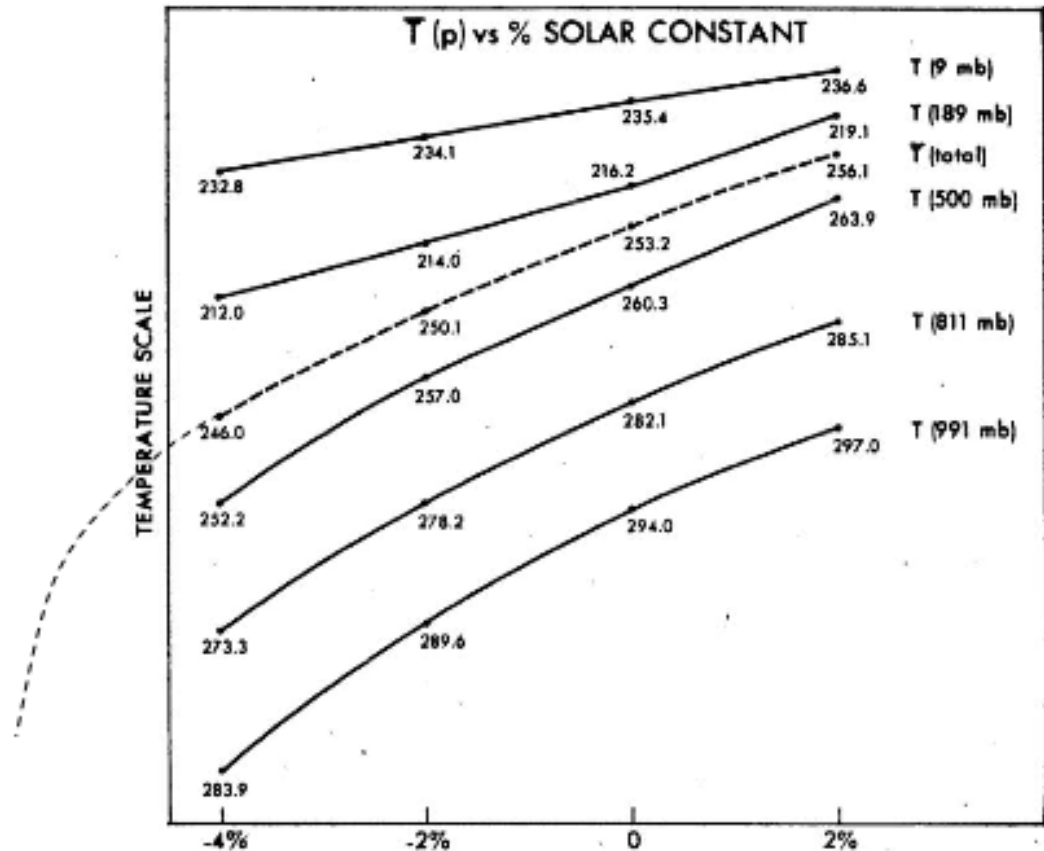
(1K per % of  $Q$ )



## Climate sensitivity to insolation in a General Circulation Model (GCM)

"As stated in the Introduction, it is not, however, reasonable to conclude that the present results are more reliable than the results from the one-dimensional studies mentioned above simply because our model treats the effect of transport explicitly rather than by parameterization."\*

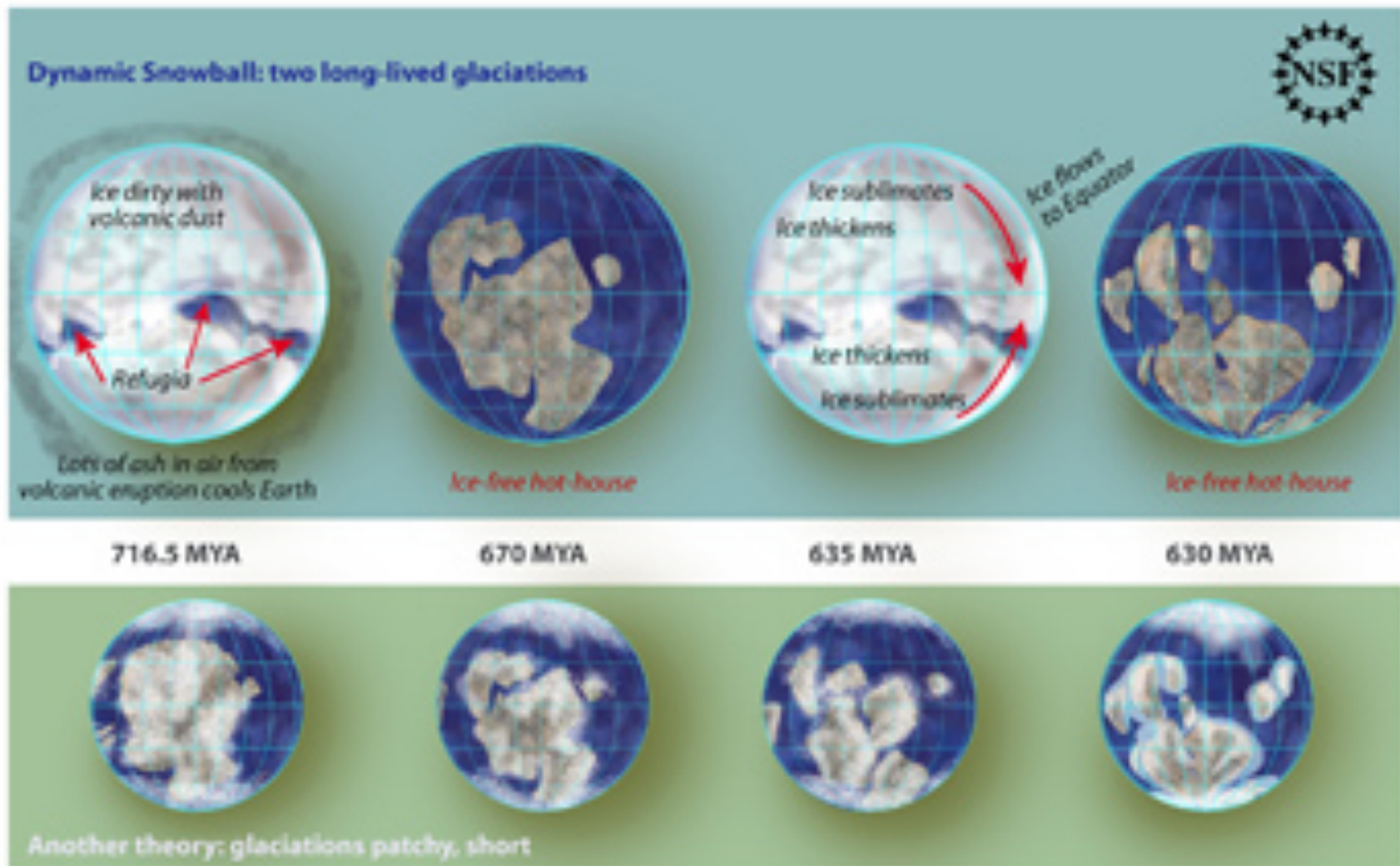
"Nevertheless, it seems to be significant that both the one-dimensional and three-dimensional models yields qualitatively similar results in many respects."\*



*Area-mean temperatures for various model levels, as well as a mass-weighted mean temperature for the total model atmosphere. Based on 4 GCM runs: control, -4%, -2% and +4%. Units are in degrees K.*

\* From Wetherald and Manabe, 1975, *J. Atmos. Sci.*, **32**, 2044–2059.

# Snowball Earth — Erstwhile a “theory”; now a “fact”?



[https://www.nsf.gov/news/news\\_images.jsp?cntn\\_id=116410&org=NSF](https://www.nsf.gov/news/news_images.jsp?cntn_id=116410&org=NSF)

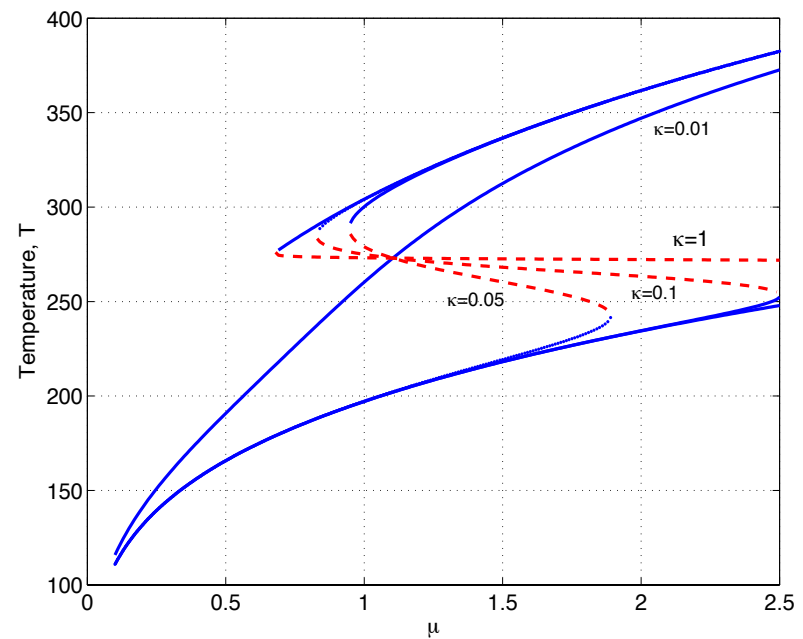
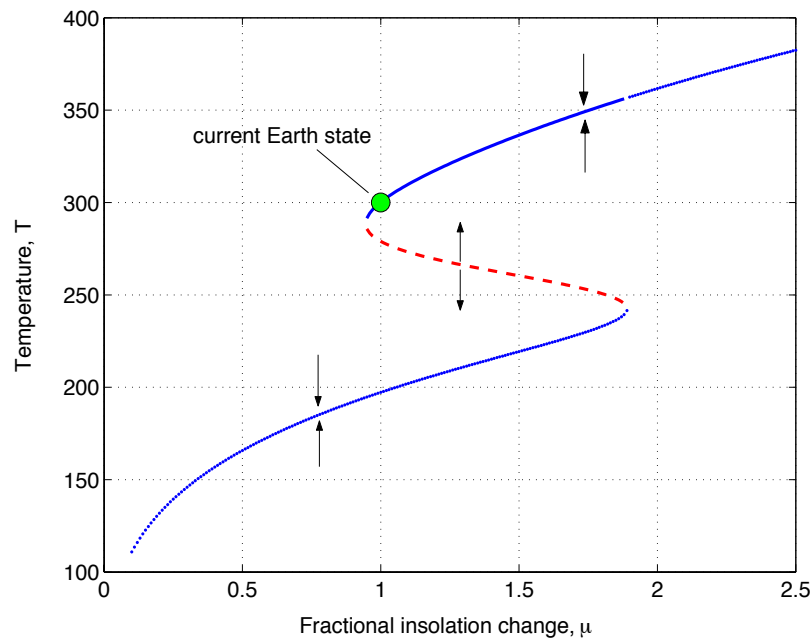
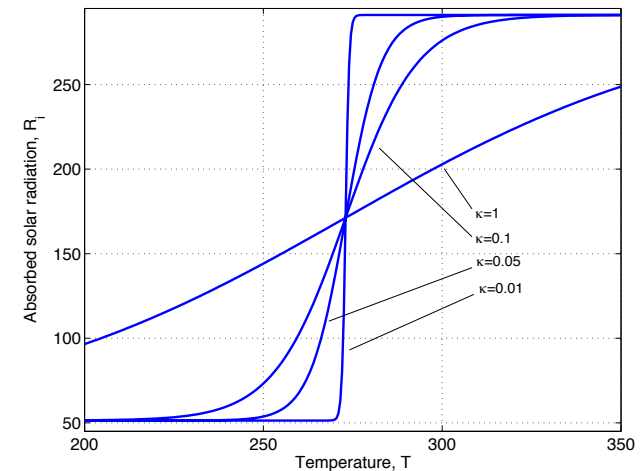
# Distance to “tipping points”?

Slightly modified 0-D EBM (Zaliapin & Ghil, *NPG*, 2010)

$$c\dot{T} = \mu Q_0(1 - \alpha(T))\sigma T^4[1 - m \tanh((T/T_0)^6)]$$

$$\alpha(T; \kappa) = c_1 + c_2 \frac{1 - \tanh[\kappa(T - T_c)]}{2}$$

$T_c$  is the ice-margin temperature,  
while  $\kappa$  is an extra “Budygo-vs.-Sellers” parameter



## Double-well potential in 2-D

1-D EBM of Budyko-Sellers-North, cf. Held & Suarez (*Tellus*, 1974); North *et al.* (*JAS*, 1979).

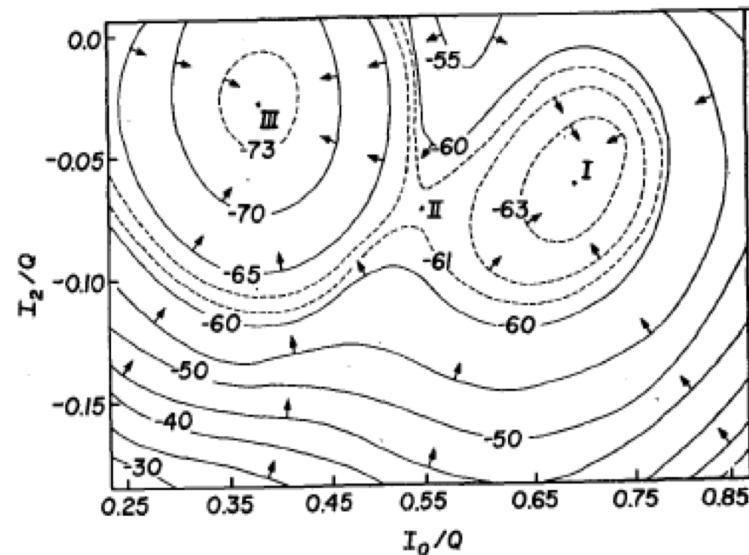
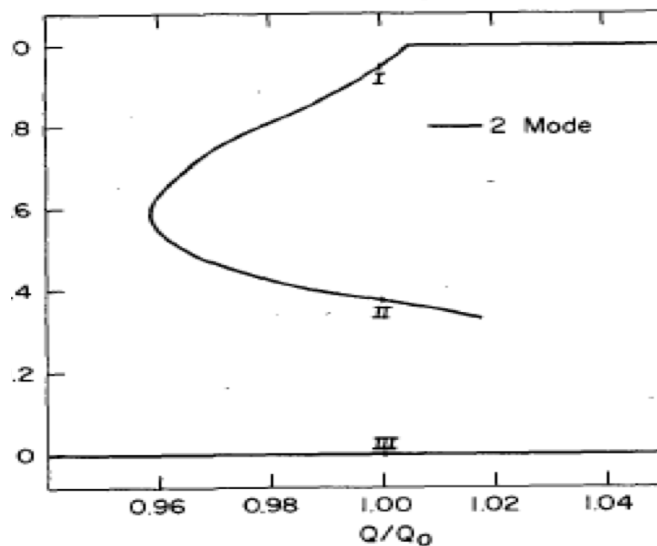
Taking  $x = \sin(\text{latitude})$  and  $k(x, T) = k_0$ ,  
We get the semi-linear parabolic PDE

$$CT_t = [k_0(1 - x^2)T_x]_x + Q(x)[1 - \alpha(T)] - I(T)$$

which yields the variational principle:

$$F\{T(x)\} = \int \left\{ \frac{1}{2}k_0(1 - x^2)T_x^2 - Q(x)A(T) + R(T) \right\} dx, \text{ where}$$

$$A(T) = \int^T [1 - \alpha(T)] dT, \text{ and } R(T) = \int^T I(T) dT.$$



# Concluding remarks, I


- ◆ Tipping points and bifurcations: multiple equilibria and rapid transitions between them.
- ◆ Prediction of the transitions? To follow.
- ◆ Transitions between more general types of behavior — limit cycles, strange attractors — likewise to follow.

# Some general references

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**Reserve slides**

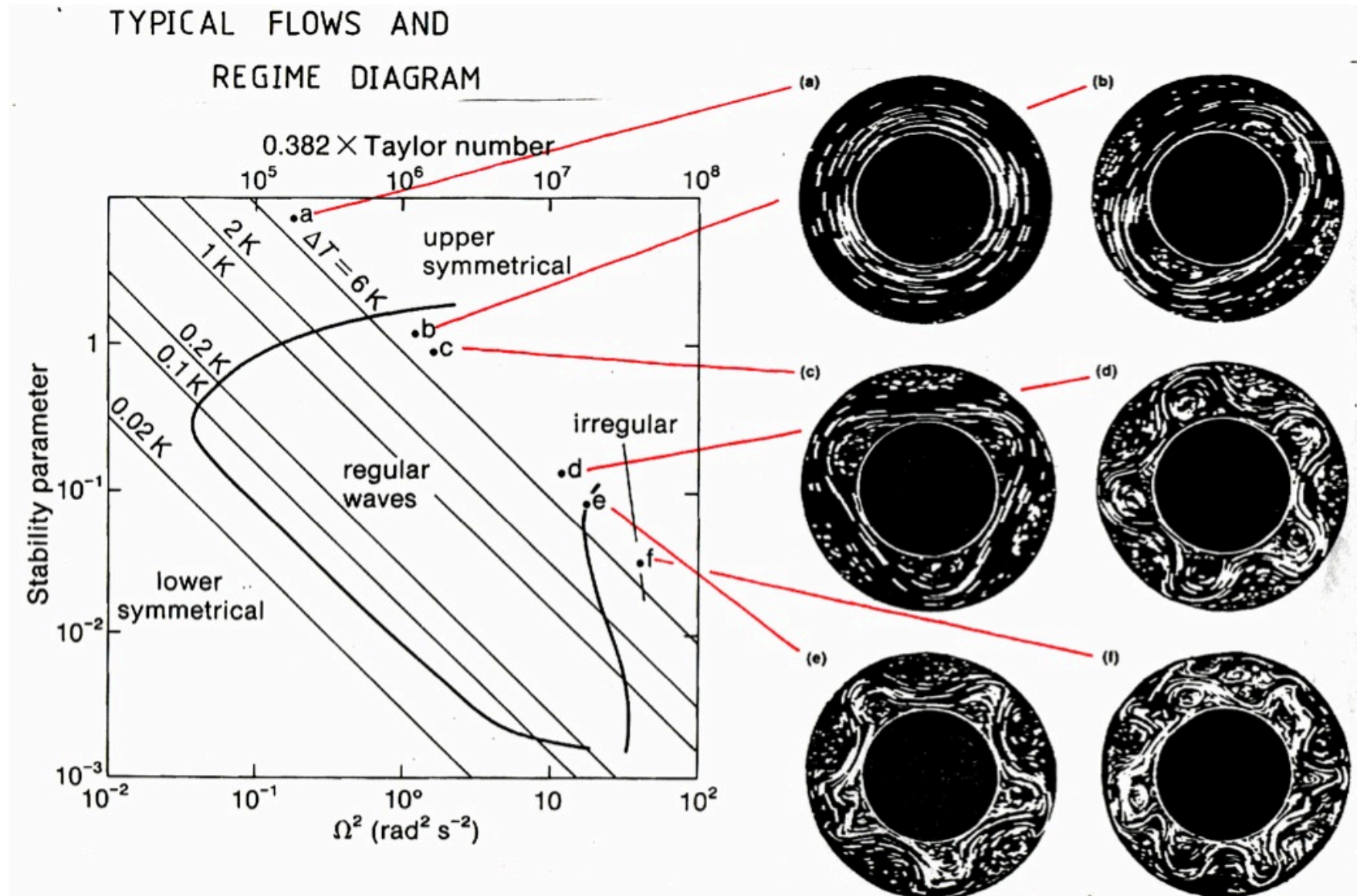
# Motivation

- There's a lot of talk about “**tipping points**.”
- It sounds **threatening**, like falling off a cliff: that's why **we care!**
- But what are they, and what do **we know** about them?
- Here's a **disambiguation page** (cf. Wikipedia), first.
- **Sociology**: “the moment of critical mass, the threshold, the boiling point” (Gladwell, 2000); a previously rare phenomenon becomes rapidly and dramatically more common.
- **Physics**: the point at which a system changes from a stable equilibrium into a new, qualitatively dissimilar equilibrium (throwing a switch, tilting a plank, boiling water, etc.).
- **Climatology**: “A climate tipping point is a **somewhat ill-defined concept** [...]” — so we'll try to actually define it better.  

- **Catastrophe theory**: branch of **bifurcation theory** in the study of **dynamical systems**; here, a tipping point is “a parameter value at which the set of equilibria abruptly change.” → **Let's see!**

M. Gladwell (2000) *The Tipping Point: How Little Things Can Make a Big Difference*.

T. M. Lenton *et al.* (2008) Tipping elements in the Earth's climate system, *PNAS*, v. **105**.

# Rotating Convection: An Illustration



# Energy Balance Models (EBMs)

## Budyko, Sellers and Held-Suarez-North

Table 10.1. Comparison of Budyko's and Sellers' models.

Heat Flux	Budyko	Sellers
$R_i = Q(1 - \alpha(T))$ Absorbed solar radiation, as a function of ice-albedo feedback	Step-function albedo $\alpha = \begin{cases} \alpha_M, & T < T_s, \\ \alpha_m, & T \geq T_s, \end{cases}$ $\alpha_M > \alpha_m,$ $T_\ell \leq T_s \leq T_u$	Ramp-function albedo $\alpha = \begin{cases} \alpha_M, & T < T_\ell \\ \alpha_M - \frac{T - T_\ell}{T_u - T_\ell} (\alpha_M - \alpha_m), & T_\ell \leq T < T_u \\ \alpha_m, & T \geq T_u \end{cases}$
$R_o$ Outgoing IR radiation	Linear, empirical $A + BT$	Stefan-Boltzmann law with greenhouse effect $\sigma T^4 \{1 - m \tanh(T^6/T_0^6)\}$
$\nabla \cdot F$ Horizontal flux divergence	Newtonian cooling $\kappa(T(\phi) - \bar{T})$	Eddy-diffusive $\nabla \cdot (k(\phi) \nabla T(\phi))$

2<sup>nd</sup> column:  
Budyko (1968, 1969)

3<sup>rd</sup> column:  
Sellers (1969)

In red:  
the "mixed" version of  
Held & Suarez (1974)  
and North (1975a, b)